



Calhoun: The NPS Institutional Archive

Faculty and Researcher Publications

Faculty and Researcher Publications

2007

Monotonicity of a key function arised in studies of nematic liquid crystal polymers

Zhou, Hong

Monotonicity of a key function arised in studies of nematic liquid crystal polymers (with H. Wang), Abstract and Applied Analysis, Vol. 2007, Article ID 76209, 7 pages, 2007.



Calhoun is a project of the Dudley Knox Library at NPS, furthering the precepts and goals of open government and government transparency. All information contained herein has been approved for release by the NPS Public Affairs Officer.

Dudley Knox Library / Naval Postgraduate School
411 Dyer Road / 1 University Circle
Monterey, California USA 93943

<http://www.nps.edu/library>

Research Article

Monotonicity of a Key Function Arised in Studies of Nematic Liquid Crystal Polymers

Hongyun Wang and Hong Zhou

Received 13 April 2007; Accepted 11 July 2007

Recommended by Nobuyuki Kenmochi

We revisit a key function arised in studies of nematic liquid crystal polymers. Previously, it was conjectured that the function is strictly decreasing and the conjecture was numerically confirmed. Here we prove the conjecture analytically. More specifically, we write the derivative of the function into two parts and prove that each part is strictly negative.

Copyright © 2007 H. Wang and H. Zhou. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.

1. Introduction

In our previous study of steady states and dynamics of two-dimensional nematic liquid crystal polymers under an imposed weak shear [1], we carried out multiscale asymptotic expansions [2, 3] to reveal the slow time evolution of the polymer orientation distribution. In the two-dimensional case, the orientation of a polymer rod is represented by angle θ [4–6]. Let $\rho(\theta, t)$ denote the probability density of the orientation of the polymer rod ensemble. Our multiscale asymptotic analysis on the Smoluchowski equation which governs the evolution of the probability density [1] demonstrates that, to the leading term, the probability density $\rho(\theta, t)$ behaves like a travelling wave with nonuniform velocity

$$\rho(\theta, t) = \rho^{(0)}(\theta - \alpha(t)) + \dots \quad (1.1)$$

Here $\rho^{(0)}(\theta)$ is a given function and the phase angle $\alpha(t)$ evolves according to

$$\frac{d\alpha(t)}{dt} = \epsilon c_1 (\sin^2(\alpha(t)) - c_2), \quad (1.2)$$

where ϵ is the Peclet number. For a weak shear, the Peclet number is small. c_1 is a positive constant independent of ϵ and α . The expression of c_2 is given by

$$c_2 = \frac{h(r)}{2g(r)}. \quad (1.3)$$

From the evolution equation (1.2), we see that the sign of c_2 determines the behavior of polymer orientation distribution: for $0 < c_2 < 1$, the polymer orientation distribution will converge to a steady state while for $c_2 < 0$, the polymer orientation distribution will keep rotating (tumbling). Function $g(r)$ is defined as

$$g(r) \equiv \left(\frac{1}{2\pi} \int_0^{2\pi} \exp(r \cos 2\theta) d\theta \right)^2. \quad (1.4)$$

Function $h(r)$ is defined as

$$h(r) \equiv \frac{1}{\langle \cos 2\theta \rangle} (1 - g(r)(1 - \langle \cos 2\theta \rangle)). \quad (1.5)$$

Here $\langle \cos 2\theta \rangle$ denotes the average of random variable $\cos 2\theta$ defined as

$$\langle \cos 2\theta \rangle \equiv \int_0^{2\pi} \cos 2\theta \rho(\theta, r) d\theta, \quad (1.6)$$

where the probability density used in the average is given by

$$\rho(\theta, r) \equiv \frac{\exp(r \cos 2\theta) d\theta}{\int_0^{2\pi} \exp(r \cos 2\theta) d\theta}. \quad (1.7)$$

It is clear that $\langle \cos 2\theta \rangle$ is a function of r . But for conciseness, we will simply write $\langle \cos 2\theta \rangle$ instead of writing $\langle \cos 2\theta \rangle(r)$. In the discussion here, r is treated as an independent variable for mathematical convenience. From physical considerations, a more meaningful quantity is, U , the normalized polymer concentration. In [1], we showed that r is a strictly increasing function of U , so there is a one-to-one correspondence between U and r .

Since function $g(r)$ is always positive, the sign of c_2 is completely determined by function $h(r)$. In [1], by doing expansions at $r = 0$ and at $r = \infty$, we showed

$$\begin{aligned} h(0) &= 1, \\ h'(0) &= -1, \\ h(r) &= \frac{-1}{4\pi r^2} \exp(2r) + \cdots \quad \text{as } r \rightarrow +\infty. \end{aligned} \quad (1.8)$$

In [1], we made a conjecture that $h(r)$ is a strictly decreasing function of r for $r > 0$. Based on the conjecture, we concluded that there is a threshold r_0 such that $h(r) > 0$ for $r < r_0$ and $h(r) < 0$ for $r > r_0$. Because of the one-to-one correspondence between U and r , the conclusion in r leads to the conclusion that there is a threshold U_0 for the normalized polymer concentration: for $U < U_0$, the polymer orientation distribution will converge to a steady state while for $U > U_0$, the polymer orientation distribution will keep rotating (tumbling).

2. Proof of the conjecture

In [1], the conjecture that $h(r)$ is a strictly decreasing function of r for $r > 0$ is only numerically confirmed. In this study, we prove this conjecture analytically. We start by writing function $h(r)$ in a slightly different form:

$$h(r) = \frac{-(g(r) - 1)}{\langle \cos 2\theta \rangle} + g(r). \quad (2.1)$$

We calculate the derivative of $\rho(\theta, r)$:

$$\begin{aligned} \frac{d}{dr}\rho(\theta, r) &= \frac{\cos 2\theta \exp(r \cos 2\theta)}{\int_0^{2\pi} \exp(r \cos 2\theta) d\theta} \\ &\quad - \frac{\exp(r \cos 2\theta) \int_0^{2\pi} \cos 2\theta \exp(r \cos 2\theta) d\theta}{\left(\int_0^{2\pi} \exp(r \cos 2\theta) d\theta\right)^2} \\ &= (\cos 2\theta - \langle \cos 2\theta \rangle) \rho(\theta, r). \end{aligned} \quad (2.2)$$

Using this result, we compute the derivative of $\langle \cos 2\theta \rangle$:

$$\begin{aligned} \frac{d}{dr}\langle \cos 2\theta \rangle &= \langle \cos 2\theta (\cos 2\theta - \langle \cos 2\theta \rangle) \rangle \\ &= \langle \cos^2 2\theta \rangle - \langle \cos 2\theta \rangle^2 \\ &= \text{var}(\cos 2\theta) > 0, \end{aligned} \quad (2.3)$$

where $\text{var}(\cos 2\theta)$ denotes the variance of random variable $\cos 2\theta$ and is always positive for finite values of r . The derivative of function $g(r)$ is given by

$$\begin{aligned} \frac{d}{dr}g(r) &= \frac{2}{4\pi^2} \int_0^{2\pi} \exp(r \cos 2\theta) d\theta \int_0^{2\pi} \cos 2\theta \exp(r \cos 2\theta) d\theta \\ &= 2 \left(\frac{1}{2\pi} \int_0^{2\pi} \exp(r \cos 2\theta) d\theta \right)^2 \frac{\int_0^{2\pi} \cos 2\theta \exp(r \cos 2\theta) d\theta}{\int_0^{2\pi} \exp(r \cos 2\theta) d\theta} \\ &= 2g(r) \langle \cos 2\theta \rangle. \end{aligned} \quad (2.4)$$

Differentiating both sides of (2.1) with respect to r and using (2.3) and (2.4) yield

$$\begin{aligned} \frac{d}{dr}h(r) &= -(g(r) - 1) \frac{d}{dr} \left(\frac{1}{\langle \cos 2\theta \rangle} \right) - \frac{dg(r)}{dr} \frac{1}{\langle \cos 2\theta \rangle} + \frac{dg(r)}{dr} \\ &= (g(r) - 1) \frac{\text{var}(\cos 2\theta)}{\langle \cos 2\theta \rangle^2} - 2g(r) + 2g(r) \langle \cos 2\theta \rangle \\ &= (g(r) - 1 - 2g(r) \langle \cos 2\theta \rangle^2) \frac{\text{var}(\cos 2\theta)}{\langle \cos 2\theta \rangle^2} \\ &\quad + 2g(r) (\text{var}(\cos 2\theta) - 1 + \langle \cos 2\theta \rangle) \\ &\equiv Q_1(r) \frac{\text{var}(\cos 2\theta)}{\langle \cos 2\theta \rangle^2} + 2g(r) Q_2(r). \end{aligned} \quad (2.5)$$

4 Abstract and Applied Analysis

In this study, our goal is to prove $dh(r)/dr < 0$ for $r > 0$. Since in (2.5), both $g(r)$ and $\text{var}(\cos 2\theta)/\langle \cos 2\theta \rangle^2$ are always positive, to achieve our goal, we only need to prove the two theorems below.

THEOREM 2.1.

$$Q_1(r) \equiv g(r) - 1 - 2g(r)\langle \cos 2\theta \rangle^2 < 0 \quad \text{for } r > 0. \quad (2.6)$$

THEOREM 2.2.

$$Q_2(r) \equiv \text{var}(\cos 2\theta) - 1 + \langle \cos 2\theta \rangle < 0 \quad \text{for } r > 0. \quad (2.7)$$

Remark 2.3. The conclusion that function $h(r)$ is strictly decreasing for $r > 0$ follows directly from these two theorems.

To facilitate the proofs of these two theorems, we introduce a functional. Specifically, the functional maps a function of θ to a function of r defined as [5]

$$F[q(\theta)](r) \equiv \int_0^{2\pi} q(\theta) \exp(r \cos 2\theta) d\theta. \quad (2.8)$$

The derivative of $F[q(\theta)](r)$ is given by

$$\begin{aligned} \frac{d}{dr} F[q(\theta)](r) &= \int_0^{2\pi} q(\theta) \cos 2\theta \exp(r \cos 2\theta) d\theta \\ &= F[q(\theta) \cos 2\theta](r). \end{aligned} \quad (2.9)$$

With the notation of functional, $\langle \cos 2\theta \rangle$ and $g(r)$ can be written as

$$\begin{aligned} \langle \cos 2\theta \rangle &= \frac{F[\cos 2\theta](r)}{F[1](r)}, \\ g(r) &= \frac{F[1]^2(r)}{4\pi^2}. \end{aligned} \quad (2.10)$$

Proof of Theorem 2.1.

Step 1. Using the notation of functional, we express function $Q_1(r)$ as

$$Q_1(r) = \frac{1}{4\pi^2} (F[1]^2(r) - 4\pi^2 - 2F[\cos 2\theta]^2(r)). \quad (2.11)$$

It is straightforward to verify that $Q_1(0) = 0$. Thus, to prove $Q_1(r) < 0$ for $r > 0$, we only need to prove $dQ_1(r)/dr < 0$ for $r > 0$. Taking the derivative of $Q_1(r)$ and with the help of (2.9), we get

$$\frac{d}{dr} Q_1(r) = \frac{2F[\cos 2\theta](r)}{4\pi^2} (F[1](r) - 2F[\cos^2 2\theta](r)). \quad (2.12)$$

Using the fact that the functional $F[\bullet]$ is linear and using the identity $\cos^2 2\theta = 1 - \sin^2 2\theta$, we obtain

$$\frac{d}{dr} Q_1(r) = \frac{2F[\cos 2\theta](r)}{4\pi^2} (2F[\sin^2 2\theta](r) - F[1](r)). \quad (2.13)$$

Since $F[\cos 2\theta](0) = 0$ and $dF[\cos 2\theta](r)/dr = F[\cos^2 2\theta](r) > 0$, we have $F[\cos 2\theta](r) > 0$ for $r > 0$. As a result, to prove $dQ_1(r)/dr < 0$ for $r > 0$, we only need to prove $2F[\sin^2 2\theta](r) - F[1](r) < 0$ for $r > 0$.

Step 2. We now prove $2F[\sin^2 2\theta](r) - F[1](r) < 0$ for $r > 0$.

Using the Taylor expansion

$$\exp(r \cos 2\theta) = \sum_{n=0}^{\infty} \frac{r^n \cos^n 2\theta}{n!}, \quad (2.14)$$

and using the fact that $\int_0^{2\pi} \cos^{2n+1} 2\theta d\theta = \int_0^{2\pi} \sin^2 2\theta \cos^{2n+1} 2\theta d\theta = 0$ to get rid of terms of odd powers in the sum, we write the expression $2F[\sin^2 2\theta](r) - F[1](r)$ as

$$\begin{aligned} 2F[\sin^2 2\theta](r) - F[1](r) \\ = \sum_{n=0}^{\infty} \frac{r^{2n}}{(2n)!} \int_0^{2\pi} (2\sin^2 2\theta \cos^{2n} 2\theta - \cos^{2n} 2\theta) d\theta. \end{aligned} \quad (2.15)$$

Applying integration by parts, we obtain

$$\begin{aligned} \int_0^{2\pi} 2\sin^2 2\theta \cos^{2n} 2\theta d\theta &= - \int_0^{2\pi} \sin 2\theta \frac{d(\cos^{2n+1} 2\theta)}{2n+1} \\ &= \frac{2}{2n+1} \int_0^{2\pi} \cos^{2n+2} 2\theta d\theta \\ &= \begin{cases} \int_0^{2\pi} \cos^{2n} 2\theta d\theta, & n = 0 \\ < \frac{2}{2n+1} \int_0^{2\pi} \cos^{2n} 2\theta d\theta, & n > 1. \end{cases} \end{aligned} \quad (2.16)$$

Substituting (2.16) into (2.15), we arrive at

$$\begin{aligned} 2F[\sin^2 2\theta](r) - F[1](r) \\ < - \sum_{n=1}^{\infty} \frac{r^{2n}}{(2n)!} \left(\frac{2n-1}{2n+1} \right) \int_0^{2\pi} \cos^{2n} 2\theta d\theta < 0. \end{aligned} \quad (2.17)$$

This completes the proof of Theorem 2.1. □

Proof of Theorem 2.2.

Step 1. Using the definition $\text{var}(\cos 2\theta) = \langle \cos^2 2\theta \rangle - \langle \cos 2\theta \rangle^2$, the relation $\langle \cos^2 2\theta \rangle - 1 = -\langle \sin^2 2\theta \rangle$, and the notation of functional, we write function $Q_2(r)$ as

$$\begin{aligned} Q_2(r) &= -\langle \sin^2 2\theta \rangle - \langle \cos 2\theta \rangle^2 + \langle \cos 2\theta \rangle \\ &= -\frac{F[\sin^2 2\theta](r)}{F[1](r)} - \frac{F[\cos 2\theta]^2(r)}{F[1]^2(r)} + \frac{F[\cos 2\theta](r)}{F[1](r)}. \end{aligned} \quad (2.18)$$

Applying integration by parts to $F[\sin^2 2\theta](r)$, we have

$$\begin{aligned}
 F[\sin^2 2\theta](r) &= \int_0^{2\pi} \sin^2 2\theta \exp(r \cos 2\theta) d\theta \\
 &= \frac{-1}{2r} \int_0^{2\pi} \sin 2\theta d(\exp(r \cos 2\theta)) \\
 &= \frac{1}{r} \int_0^{2\pi} \cos 2\theta \exp(r \cos 2\theta) d\theta \\
 &= \frac{1}{r} F[\cos 2\theta](r).
 \end{aligned} \tag{2.19}$$

Substituting (2.19) into (2.18), using the fact that the functional $F[\bullet]$ is linear and using the relation $\cos 2\theta = 1 - 2\sin^2 \theta$, we obtain

$$\begin{aligned}
 Q_2(r) &= \frac{F[\cos 2\theta](r)}{F[1]^2(r)} \left(\left(1 - \frac{1}{r}\right) F[1](r) - F[\cos 2\theta](r) \right) \\
 &= \frac{1}{r} \cdot \frac{F[\cos 2\theta](r)}{F[1]^2(r)} (2F[r\sin^2 \theta](r) - F[1](r)).
 \end{aligned} \tag{2.20}$$

In the proof of Theorem 2.1, we have shown that $F[\cos 2\theta](r) > 0$ for $r > 0$. It follows that to prove $Q_2(r) < 0$ for $r > 0$, we only need to show $2F[r\sin^2 \theta](r) - F[1](r) < 0$ for $r > 0$.

Step 2. We now prove $2F[r\sin^2 \theta](r) - F[1](r) < 0$ for $r > 0$.

Using $\cos 2\theta = 2\cos^2 \theta - 1$, we rewrite and expand $\exp(r \cos 2\theta)$ as

$$\begin{aligned}
 \exp(r \cos 2\theta) &= \exp(-r) \exp(2r \cos^2 \theta) \\
 &= \exp(-r) \sum_{n=0}^{\infty} \frac{(2r)^n \cos^{2n} \theta}{n!}.
 \end{aligned} \tag{2.21}$$

Using this result, we write $2F[r\sin^2 \theta](r) - F[1](r)$ as

$$\begin{aligned}
 &2F[r\sin^2 \theta](r) - F[1](r) \\
 &= \exp(-r) \left[\sum_{n=0}^{\infty} \frac{(2r)^{n+1}}{n!} \int_0^{2\pi} \sin^2 \theta \cos^{2n} \theta d\theta - \sum_{n=0}^{\infty} \frac{(2r)^n}{n!} \int_0^{2\pi} \cos^{2n} \theta d\theta \right] \\
 &= \exp(-r) \left[\sum_{n=0}^{\infty} \frac{(2r)^{n+1}}{n!} \int_0^{2\pi} \sin^2 \theta \cos^{2n} \theta d\theta - 2\pi - \sum_{m=0}^{\infty} \frac{(2r)^{m+1}}{(m+1)!} \int_0^{2\pi} \cos^{2(m+1)} \theta d\theta \right] \\
 &= \exp(-r) \left[\sum_{n=0}^{\infty} \frac{(2r)^{n+1}}{n!} \int_0^{2\pi} \left(\sin^2 \theta \cos^{2n} \theta - \frac{\cos^{2n+2} \theta}{n+1} \right) d\theta - 2\pi \right].
 \end{aligned} \tag{2.22}$$

Integrating by parts, we obtain

$$\begin{aligned}
 \int_0^{2\pi} \sin^2 \theta \cos^{2n} \theta d\theta &= - \int_0^{2\pi} \sin \theta \frac{d(\cos^{2n+1} \theta)}{2n+1} \\
 &= \frac{1}{2n+1} \int_0^{2\pi} \cos^{2n+2} \theta d\theta.
 \end{aligned} \tag{2.23}$$

Substituting (2.23) into (2.22), we arrive at

$$2F[r\sin^2\theta](r) - F[1](r) = \exp(-r) \left[- \sum_{n=1}^{\infty} \frac{(2r)^{n+1}}{n!} \frac{n}{(2n+1)(n+1)} \int_0^{2\pi} \cos^{2n+2}\theta d\theta - 2\pi \right] < 0 \quad \text{for } r > 0. \quad (2.24)$$

This completes the proof of Theorem 2.2. \square

3. Concluding remarks

We provide a rigorous proof on the monotonicity of a key function in the study of nematic liquid crystal polymers. This monotonicity has enabled us to establish analytically the tumbling region of shear-driven nematic liquid crystal polymers [1]. This paper serves as a theoretical advance in our studies of complex fluids.

Acknowledgment

This work was partially supported by the U.S. National Science Foundation and the Air Force Office of Scientific Research.

References

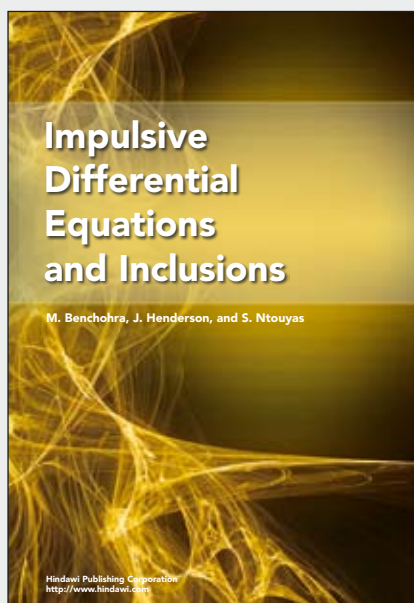
- [1] H. Zhou and H. Wang, “Steady states and dynamics of 2-D nematic polymers driven by an imposed weak shear,” *Communications in Mathematical Sciences*, vol. 5, no. 1, pp. 113–132, 2007.
- [2] A. W. Bush, *Perturbation Methods for Engineers and Scientists*, CRC Press, Boca Raton, Fla, USA, 1992.
- [3] E. J. Hinch, *Perturbation Methods*, Cambridge Texts in Applied Mathematics, Cambridge University Press, Cambridge, UK, 1991.
- [4] C. Luo, H. Zhang, and P. Zhang, “The structure of equilibrium solutions of the one-dimensional Doi equation,” *Nonlinearity*, vol. 18, no. 1, pp. 379–389, 2005.
- [5] P. Constantin and J. Vukadinovic, “Note on the number of steady states for a two-dimensional Smoluchowski equation,” *Nonlinearity*, vol. 18, no. 1, pp. 441–443, 2005.
- [6] I. Fatkullin and V. Slastikov, “A note on the Onsager model of nematic phase transitions,” *Communications in Mathematical Sciences*, vol. 3, no. 1, pp. 21–26, 2005.

Hongyun Wang: Department of Applied Mathematics and Statistics, University of California, Santa Cruz, CA 95064, USA
Email address: hongwang@ams.ucsc.edu

Hong Zhou: Department of Applied Mathematics, Naval Postgraduate School, Monterey, CA 93943, USA
Email address: hzhou@nps.edu

Impulsive Differential Equations and Inclusions

M. Benchohra, J. Henderson, and S. Ntouyas



Functional differential equations and inclusions have assumed an increasing relevance, mainly due to the great variety of problems arising in sciences, engineering, and technology which can be described by such types of equations and inclusions. This book is devoted to this rapidly developing theory of differential equations and inclusions with special emphasis on such equations and inclusions subject to impulses. In addition, the book includes a wide account of various applications to physical problems.

The book consists of twelve chapters. Initial value problems and boundary value problems for both impulsive differential equations and differential inclusions, as well as for both impulsive functional differential equations and functional differential inclusions, or for both neutral functional differential equations and neutral functional differential

inclusions, are studied. Moreover, results on impulsive semilinear ordinary differential inclusions and functional differential inclusions satisfying nonlocal boundary conditions are given extensive attention. Positive solutions and multiple positive solutions for impulsive ordinary differential equations and functional differential equations, as well as impulsive differential inclusions satisfying periodic boundary conditions, are also studied.

Some of these results are extended to functional differential equations and functional differential inclusions with variable times, to semilinear functional differential equations and functional differential inclusions for operators that are nondensely defined on a Banach space, to impulsive hyperbolic differential inclusions, and to impulsive dynamic equations on time scales.

Each chapter concludes with a section devoted to notes and bibliographical remarks. This book is addressed to a wide audience of specialists such as mathematicians, engineers, biologists, and physicists.

Limited-Time
Promotional Offer.
Buy this title NOW at
20% discount plus
Free Shipping.

CMIA Book Series, Volume 2, ISBN: 977-5945-50-X

Please visit <http://www.hindawi.com/cmia.2.html> for more information about the book.

To place an order while taking advantage of our current promotional offer, please contact books.orders@hindawi.com